

Statistical mechanics in the nanoscale: from physics to biology

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Recent review: A. Alemany, M. Ribezzi and F.R. in *Nonequilibrium Statistical Physics of Small Systems: Fluctuation Relations and Beyond*, R.Klages, W.Just, C.Jarzynski (Eds.), Wiley-VCH 2012 Work also done in collaboration with A. Crisanti and M. Picco, E. Dieterich and U. Seifert

About nonequilibrium states

In equilibrium:

 No net energy & mass transfer between system and sources

Boltzmann-Gibbs distribution

Out of equilibrium:

- Net energy & mass transfer
- Breakdown of maximum entropy principle
- Non-conservative or time-dependent forces

For a review, F. Ritort, Adv. Chem. Phys. 137 31 (2008)

WEAKLY ERGODIC OR GLASSY SYSTEMS



Slow dynamics or aging | FDT violations Intermittent events | Kinetic heterogeneities

Non-equilibrium fluctuations due to **small** cooperatively rearranging regions (CRRs)

AGING



Simulations: Structural glasses (Kob, Barrat 98) and spin glasses (Picco, Ricci-Tersenghi, F.R., 2000)



FDT violations and effective temperatures

$$k_{B}T_{eff}(\omega) = \frac{\pi\omega S(\omega)}{2\hat{\chi}''(\omega)} \begin{cases} \text{Power spectrum} \\ \text{(fluctuations)} \\ \text{Susceptibility} \\ \text{(response)} \end{cases}$$

Two-times response Two-times correlation

$$T_{eff}(t_w)R(t,t_w) = \frac{\partial C(t,t_w)}{\partial t_w}$$

Cugliandolo and Kurchan, Phys. Rev. Lett. 71, 173 (1993

Q: Is there a thermodynamic nonequilibrium temperature?

L.F.Cugliandolo, J.Kurchan and L. Peliti, Phys. Rev. E, 55, 3898 (1998)

For "numerical evidence", see A. Crisanti and F. Ritort, J. Phys. A (Math. Gen.) 36, R181 (2003)

Fluctuation-dissipation plots (FD plots)

$$\frac{1}{T_{eff}(t_w)} = \frac{\partial \chi(t,t_w)}{\partial C(t,t_w)}\Big|_{t_w \text{ fixed}}$$

L.F. Cugliandolo and J. Kurchan, J. Phys. A (Math. Gen.) **27**, 5749 (1993)

For "numerical evidence", see A. Crisanti and F. Ritort, J. Phys. A (Math. Gen.) 36, R181 (2003)



Fluctuation-dissipation plots



Main idea: glassy systems exchange energy with the environment in a process dominated by rare and intermittent events (e.g. relaxation of internal stress) and large deviations respect to the average behavior.

Strategy: The understanding of glassy behavior can be approached from the perspective of nonequilibrium thermodynamics of small systems (also called stochastic thermodynamics) and fluctuation theorems.



F. R. in *Unifying concepts in Granular Media and Glassy systems,* Ed. By A. Coniglio, A. Fierro and H. Hermann, Springer Verlag, Berlin 2004, page 129



 $\begin{bmatrix} 10^{2} \\ 0.5 \\ 10^{9} \\ 10^{2} \end{bmatrix} = \begin{bmatrix} 0.7 \\ 0.6 \\ 4 \times 10^{3} \\ 5 \times 10^{3} \\ 0.6 \\ t_{w} (sec) \end{bmatrix} = \begin{bmatrix} 0.7 \\ 0.6 \\ t_{w} (sec) \\ 0 \\ 0.55 \\ 0.60 \\ 0.65 \\ c_{I}(t_{w}, \tau) \end{bmatrix}$

Confocal microscopy in colloids:L.Cipelletti (Montpellier) and E.Weeks (USA) groups

Noise measurements in glasses (laponite and polycarbonate)S.Ciliberto group in Lyon (France)

Also in steady states (theory)

Heat distribution for a particle confined in a quadratic potential and dragged through a solvent (E.D.G. *Cohen and R. Van Zon, 2003*)



Glass transition in the random orthogonal model

A random matrix (Ising spin glass) model:

E. Marinari, G. Parisi, F. R. J. Phys. A 27, 7647 (1994)





ROM : N=64



FD plots in the random orthogonal model (ROM)

Parameters: N=64,
$$T_i = 3$$
, $T = 0.2$



Heat pdfs in the random orthogonal model (ROM)

Parameters: N=64,
$$T_i = 3$$
, $T = 0.3, 0.2, 0.1$, $t_w = 2^{10}, 2^{15}$



A. Crisanti and F. R, Europhys. Lett. 66, 253 (2004)

Can we experimentally test these ideas?

Use small systems as models where nonequilibrium fluctuations and large deviations are experimentally measurable

A NEW APPROACH TO UNDERSTAND THE EMERGENCE OF EFFECTIVE TEMPERATURES

Work done in collaboration with E. Dieterich and U. Seifert, Stuttgart

Question: Can we think of a nonequilibrium model that mimics the effect of the random thermal (Brownian) forces?

Hypothesis: In principle, any external stochastic forcing applied to the system should be indistinguishable from a thermal-like Brownian force if the **characteristic timescale** of the stochastic force is shorter than the **intrinsic relaxation time** of the system

$$< f(0)f(t) >= A \exp(-t/\tau)$$

Strategy: Use **single molecules** with **stochastic driving** to generate **nonequilibrium steady states** with effective temperatures

Minitweezers: Experimental set-up

C. Bustamante and S. B. Smith et al., US Patent,**7**, 133, 132, B2 (2006)



THE WORLD'S SHORTEST HANDLES



Pulling experiments



Hopping traces (passive mode)



Hopping traces (force feedback)



A DNA hairpin with many intermediates

Optical tweezers: B. Essevaz-Roulet, U. Bockelmann , F. Heslot F (1997) *Proc Natl Acad Sci* USA 94:11935-11940 **AFM:** M. Rief, H. Clausen-Schaumann, H.E. Gaub (1999) *Nat Struct Biol* 6:346-349 **Magnetic tweezers:** C. Danilowicz et al. (2003) *Proc Natl Acad Sci* USA 100:1694-1699



N. Forns et al., Biophys. J. **100** 1765-1774 (2011)

What happens with controlled force?



Extremely slow kinetics





Experimental realization of the Sinai model: $E(t)-E(\infty) \rightarrow a/log(t)$

To dissipate a few per cent of these values (aprox. 20kT like in the trap ensemble) would require a loading rate of

10⁻¹⁰⁰ pN/s !!!!



Take home messages....

- In glassy systems an effective temperature describes the relation between correlations and responses at low frequencies, probably due to intermittent relaxational events that are indistinguishable from Brownian-like forces.
- Small systems and stochastic thermodynamics are the ideal framework to theoretically understand and experimentally measure effective temperatures.
- Noise experiments much needed: single molecules are excellent candidate systems. The high signal-to-noise ratios of single molecule techniques have provided new tools for exploring the world of nonequilibrium phenomena



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